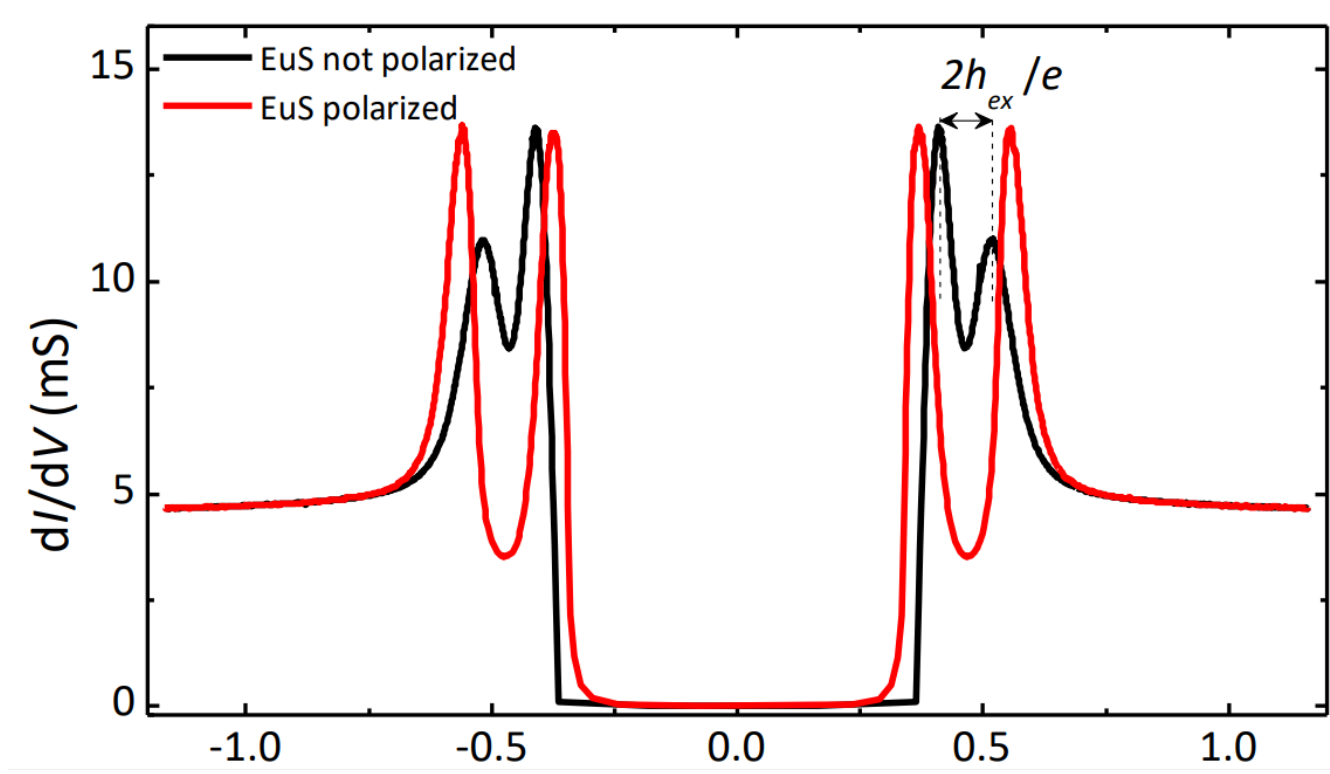


# Influence of magnons on the superconducting state in superconductor/magnet heterostructures

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## Introduction

In a thin-film ferromagnetic insulator/superconductor (FI/S) bilayers the exchange field of FIs can split the excitation spectrum of the adjacent thin-film superconductor [1]. A series of interesting phenomena have been predicted to occur in S/FI structures with spin-split DOS, such as huge thermoelectric effects [2]. By now the exchange-induced spin-splitting of the excitation spectra and DOS is studied theoretically in the mean-field approximation with respect to the exchange field.



DOS for EuS/Al bilayer. Experimental results adopted from [1].

Here we investigate the influence of electron-magnon interaction on these electronic characteristics.

## Model and method

$$\hat{H} = \hat{H}_S + \hat{H}_{FI} + \hat{H}_{ex}$$

$$\hat{H}_{FI} = \sum_q (\omega_0 + Dq^2) b_q^\dagger b_q,$$

$$\hat{H}_S = \sum_{k\sigma} \xi_k c_{k\sigma}^\dagger c_{k\sigma} - \sum_k \Delta c_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger - \sum_k \Delta^* c_{-k\downarrow} c_{k\uparrow},$$

$$\hat{H}_{ex} = -J \int d^2\rho \mathbf{S}_{FI}(\rho) \mathbf{s}_e(\rho) = \tilde{U} \sum_k (c_{k,\uparrow}^\dagger c_{k,\uparrow} - c_{k,\downarrow}^\dagger c_{k,\downarrow}) + V \sum_{k,q} (b_q c_{k,\uparrow}^\dagger c_{k-q,\downarrow} + b_q^\dagger c_{k-q,\downarrow}^\dagger c_{k,\uparrow});$$

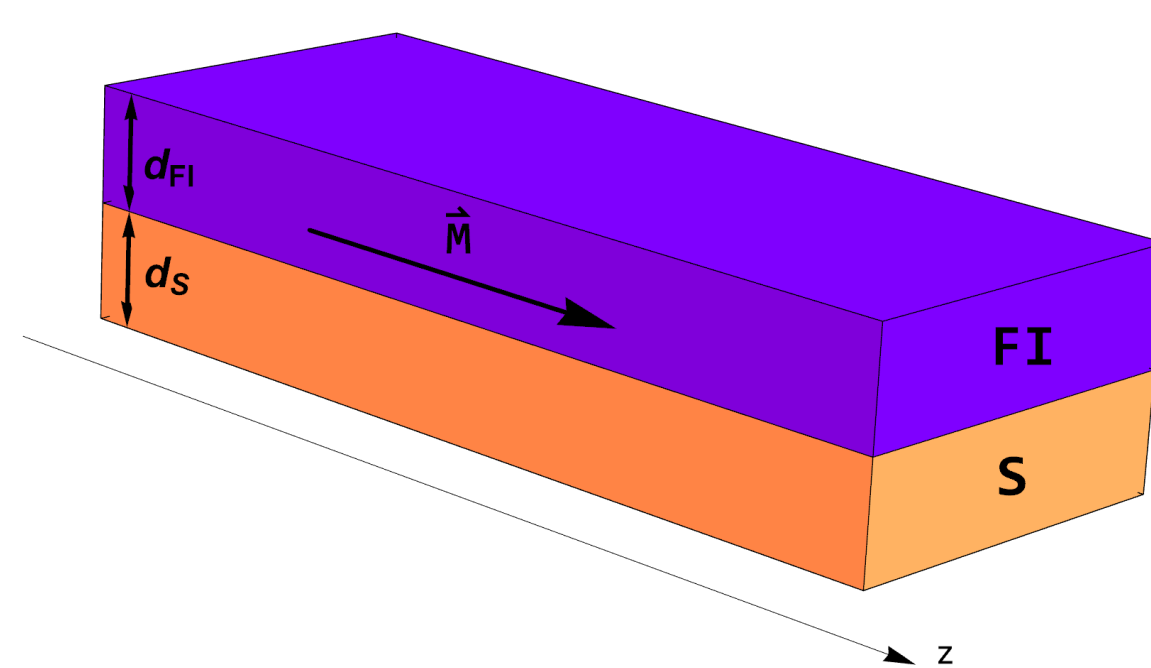
The equation on the Gor'kov Green's function:

$$(\varepsilon - \xi_k \tau_z - \sigma \tilde{U} - \Delta \tau_x - \hat{\Sigma}_{m,\sigma}) \hat{G}_{k,\sigma}(\omega) = 1,$$

$$\hat{\Sigma}_{m,\sigma} = -V^2 T \sum_{q,\Omega} B_q(\Omega) \hat{G}_{0,k-\sigma q,\sigma}(\omega - \sigma\Omega),$$

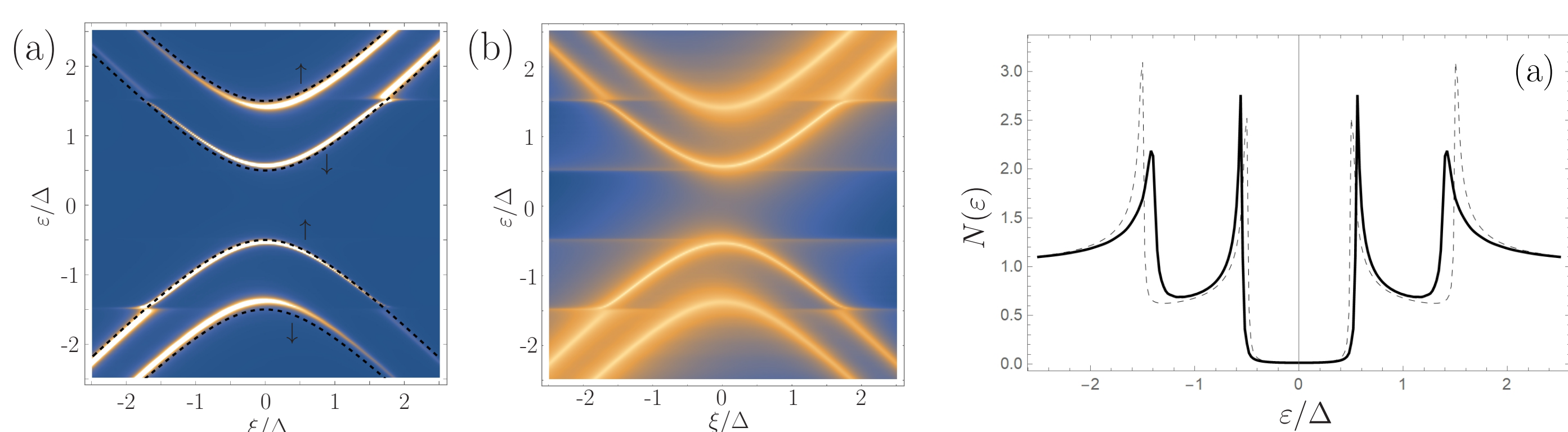
$$\tilde{U} = -\frac{J}{2|\gamma|d_S} (M_s - N_m |\gamma|), \quad B_q(\Omega) = \frac{1}{i\Omega - (\omega_0 + Dq^2)}$$

System under consideration:



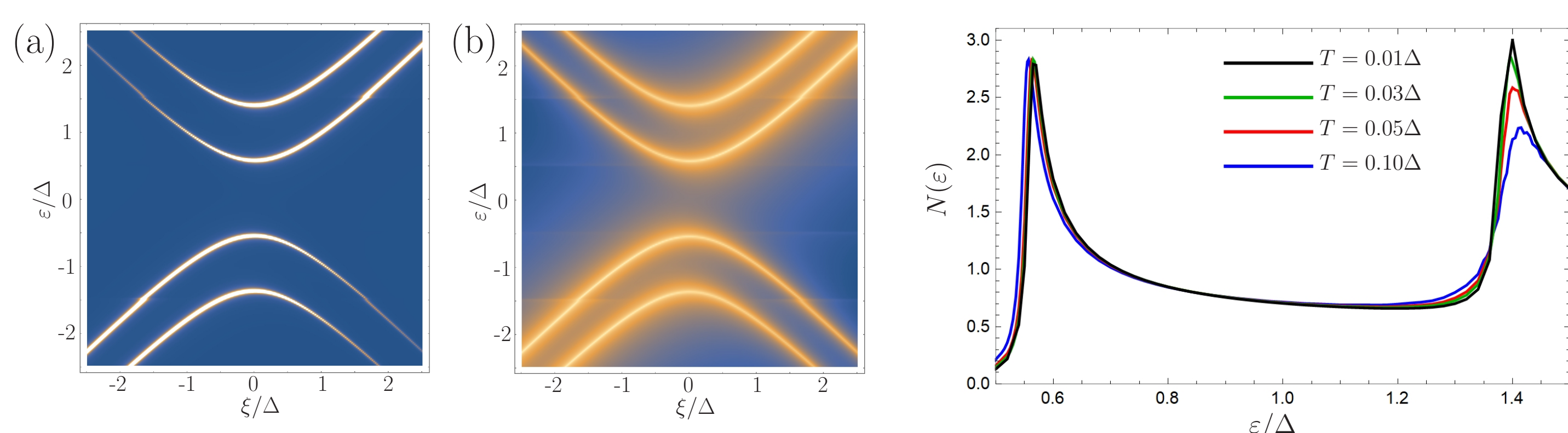
## LDOS and the dispersion law of S layer

K is the dimensionless constant quantifying the strength of the electron-magnon coupling:  $K = \frac{V^2 S_F}{4\pi\sqrt{D} d_S} = \frac{U^2 (2|\gamma|/M_s)}{4\pi d_F \Delta_0^{3/2} \sqrt{D} \xi_0}$ .



Electronic spectral function in the S layer for the both spins (a,b) and spin-up (c,d) at  $K = 0.01$ ,  $T = 0.1\Delta$ ,  $\tilde{U} = 0.5\Delta$ . Dashed lines are quasiparticle spectra at  $K=0$ .

(a) LDOS in the S layer. Dashed line is the LDOS at  $K=0$ .  
 (b) Spin-resolved LDOS.  $K = 0.01$ ,  $T = 0.1\Delta$ ,  $\tilde{U} = 0.5\Delta$ .



Electronic spectral function in the S layer for the both spins at  $K = 0.01$ ,  $T = 0.01\Delta$ ,  $\tilde{U} = 0.5\Delta$ .

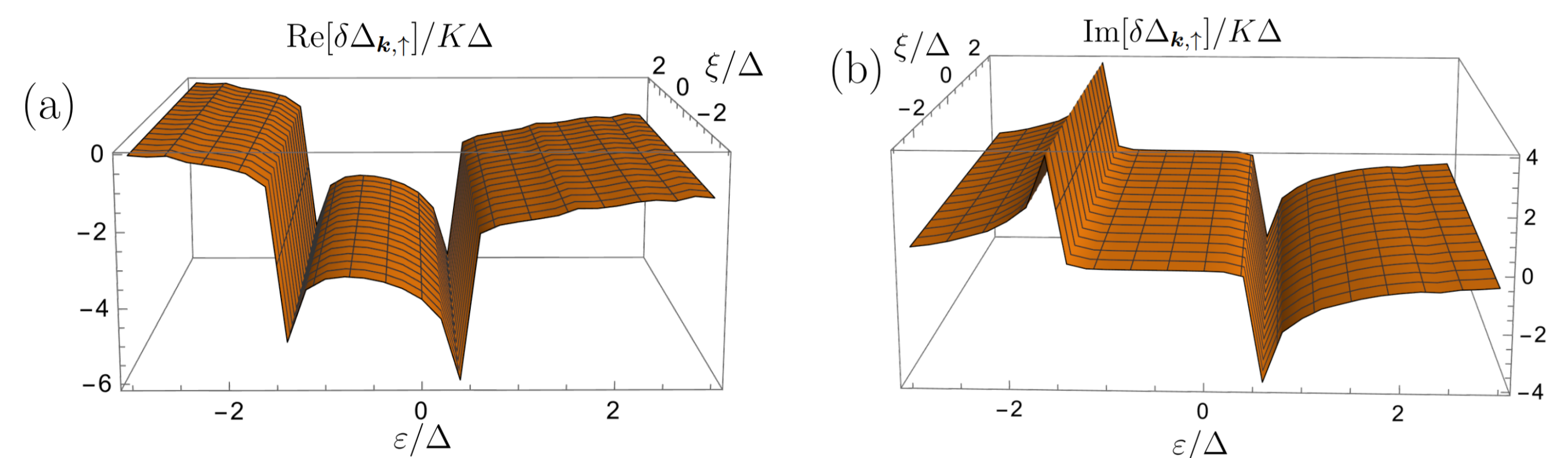
LDOS for the different temperatures.  $K = 0.01$ ,  $\tilde{U} = 0.5\Delta$ .

Electron-magnon interaction leads to the temperature-dependent reconstruction of the quasiparticle spectra and LDOS.

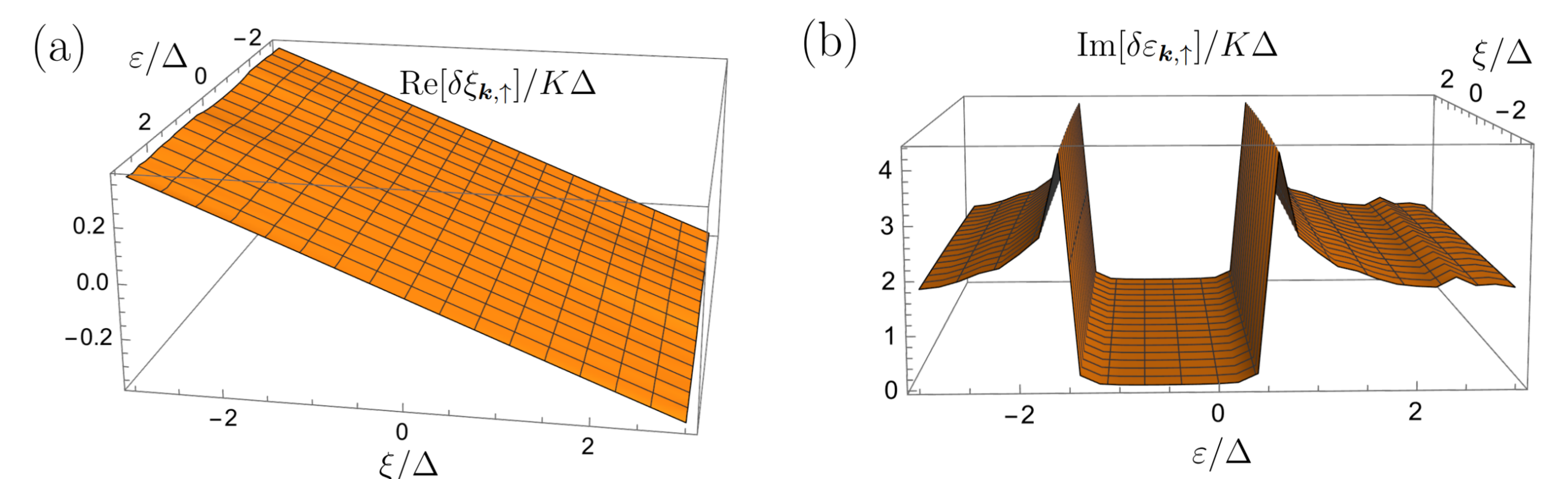
## Magnonic corrections to the Green's function

$$\hat{G}_{k,\sigma}(\varepsilon) = \frac{\tilde{\varepsilon}_{k,\sigma} + \tilde{\xi}_{k,\sigma} \tau_z + \tilde{\Delta}_{k,\sigma} \tau_x}{(\tilde{\varepsilon}_{k,\sigma})^2 - (\tilde{\xi}_{k,\sigma})^2 - (\tilde{\Delta}_{k,\sigma})^2}$$

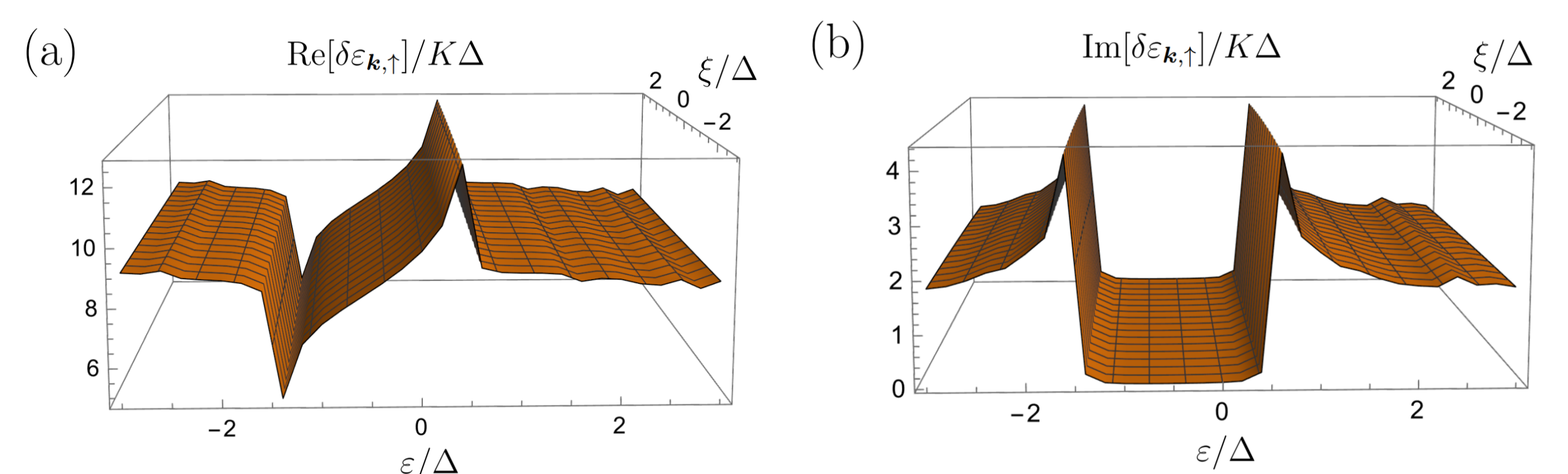
$$\tilde{\Delta}_{k,\sigma}(\varepsilon) = \Delta + \delta\Delta_{k,\sigma}(\varepsilon) = \Delta - V^2 T \sum_{q,\Omega} B_q(\Omega) \frac{\Delta}{(\varepsilon - i\sigma\Omega + \tilde{U}\sigma)^2 - \xi_{k-\sigma q}^2 - |\Delta|^2}$$



$$\tilde{\xi}_{k,\sigma}(\varepsilon) = \xi_k + \delta\xi_{k,\sigma}(\varepsilon) = \xi_k - V^2 T \sum_{q,\Omega} B_q(\Omega) \frac{\xi_{k-\sigma q}}{(\varepsilon - i\sigma\Omega + \tilde{U}\sigma)^2 - \xi_{k-\sigma q}^2 - |\Delta|^2}$$



$$\tilde{\varepsilon}_{k,\sigma}(\varepsilon) = \varepsilon - \tilde{U}\sigma + \delta\varepsilon_{k,\sigma}(\varepsilon) = \varepsilon - \tilde{U}\sigma + V^2 T \sum_{q,\Omega} B_q(\Omega) \frac{\varepsilon - i\sigma\Omega + \tilde{U}\sigma}{(\varepsilon - i\sigma\Omega + \tilde{U}\sigma)^2 - \xi_{k-\sigma q}^2 - |\Delta|^2}$$



$$\delta\varepsilon_{\downarrow}(\varepsilon) = -[\delta\varepsilon_{\uparrow}(-\varepsilon)]^*, \quad \delta\xi_{k,\downarrow}(\varepsilon) = [\delta\xi_{k,\uparrow}(-\varepsilon)]^*, \quad \delta\Delta_{\downarrow}(\varepsilon) = [\delta\Delta_{\uparrow}(-\varepsilon)]^*,$$

in Matsubaras  $\delta\Delta_{\uparrow}(\omega) = \delta\Delta_{\downarrow}(-\omega) \rightarrow \Delta_t(\omega) = -\Delta_t(-\omega)$  odd - frequency order parameter

## Conclusions

Interaction of electrons with thermal magnons in thin-film FI/S bilayers modifies the quasiparticle spectra and the DOS of the S film:

- The effective spin splitting of the coherence peaks is reduced with respect to the mean-field consideration.
- The outer spin-split coherence peaks are broadened, and the inner peaks remain intact. This type of broadening is a clear signature of the magnon-mediated spin flips and strongly differs from other mechanisms of the coherence peaks broadening, which usually influence all peaks.
- The order parameter acquires an effective odd-frequency spin-triplet component.
- The spin-split quasiparticle branches are partially mixed and reconstructed due to the magnon-mediated electron spin-flip processes.

## References

- [1] E. Strambini, V. N. Golovach, G. De Simoni, J. S. Moodera, F.S. Bergeret, F. Giazotto, Revealing the magnetic proximity effect in EuS/Al bilayers through superconducting tunneling spectroscopy, Phys. Rev. Materials **1**, 054402 (2017).
- [2] F. Sebastian Bergeret, Mikhail Silaev, Pauli Virtanen, and Tero T. Heikkila, Colloquium: Nonequilibrium effects in superconductors with a spin-splitting field, Rev. Mod. Phys. **90**, 041001 (2018).

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